

CLAS12 Techniques of Detection

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Abstract

7 In the realm of particle physics, the ability to detect the unseen plays a crucial role
8 in the field. This paper delves into the detection techniques employed at Jefferson Lab.
9 Specifically the CLAS12 detector. It provides an overview of the principles being these
10 techniques as well as use cases such as the Heavy Proton Search (HPS) and Beam Spin
11 Asymmetries. It highlights advantages and limitations associated with each technique,
12 emphasizing the importance of selecting the most appropriate method for a specific
13 detection task.

14 The study also explores the intricate mechanisms of the CLAS12 detector, a cutting-
15 edge tool in particle physics research. It discusses the detector's design, functionality,
16 and its significant contributions to the field. The paper further investigates the chal-
17 lenges faced in detecting elusive particles and the innovative solutions implemented to
18 overcome these obstacles. It underscores the importance of continuous advancements
19 in detection technology to expand our understanding of particle physics. The paper
20 serves as a comprehensive guide for researchers and students alike, offering valuable
21 insights into the complex world of particle detection at Jefferson Lab.

Keywords

- 23 – CLAS12
- 24 – Jefferson Laboratory
- 25 – Hall B

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1 Toroidal and Solenoidal Magnet

1.1 Introduction

The CLAS12 detector is a marvel of engineering designed to delve into the secrets of nuclear and hadronic interactions triggered by electrons. At its heart lies a powerful combination of magnets a six-coil toroid and a robust solenoid.

The toroidal magnet generates a strong magnetic field that extends outward, encompassing particles emitted at angles up to 35 degrees from the beam direction. Inside this toroidal field, a powerful solenoid magnet takes over, creating a 5 Tesla field that covers a central region with a polar angle range of roughly 35 to 125 degrees. This combined magnetic field configuration allows CLAS12 to efficiently detect a wide range of charged and neutral particles across a significant portion of the entire solid angle.

This configuration was suggested to enable the precise measurement of charged particles with high momenta and resolution at forward angles, while maintaining high luminosity in the detector systems. To achieve this, it's crucial to shield low-energy leptons generated in the target material via processes like Møller scattering $l^- + l^- \rightarrow l^- + l^-$ (where "l" represents a lepton). This shielding is facilitated by a robust longitudinal magnetic field that redirects the electrons towards a shielding conduit crafted from dense tungsten material, where they dissipate their energy. [1]

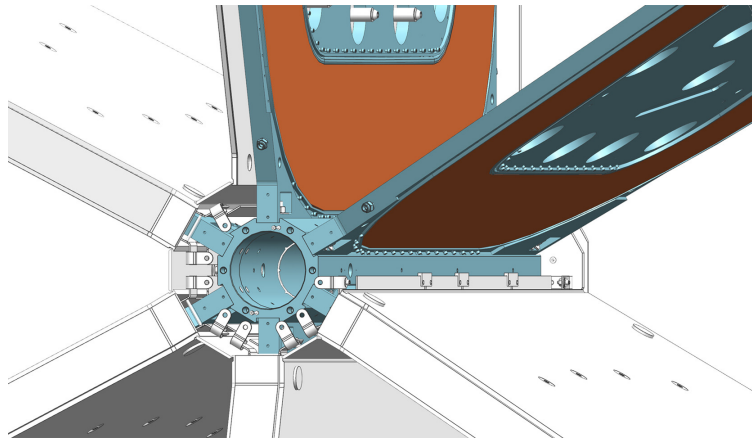


Figure 1: Torus Magnet.

67 1.2 The Torus Magnet

68 The magnetic field around the beam line is approximately toroidal, generated by six symmet-
69 rically placed magnetic coils. The sextet of coils were mounted onto a shared stainless-steel
70 cylinder within a central cold hub, providing both structural stability and geometric unifor-
71 mity to align the coils close to the magnet's core (see Fig. 1). This arrangement heightens
72 the accuracy of positioning the coil assemblies in areas where the magnetic field is expected
73 to reach its zenith.

74 Each superconducting coil consists of a two-coil (double-pancake) encased in an aluminum
75 shell. Each pancake contains 117 windings. The magnet's inductance is 2.0 H, and the stored
76 energy is 14.2 MJ. The magnet is equipped with N_2 cooled heat shields. Upon assembly and
77 cooling, the magnet promptly achieved full field.

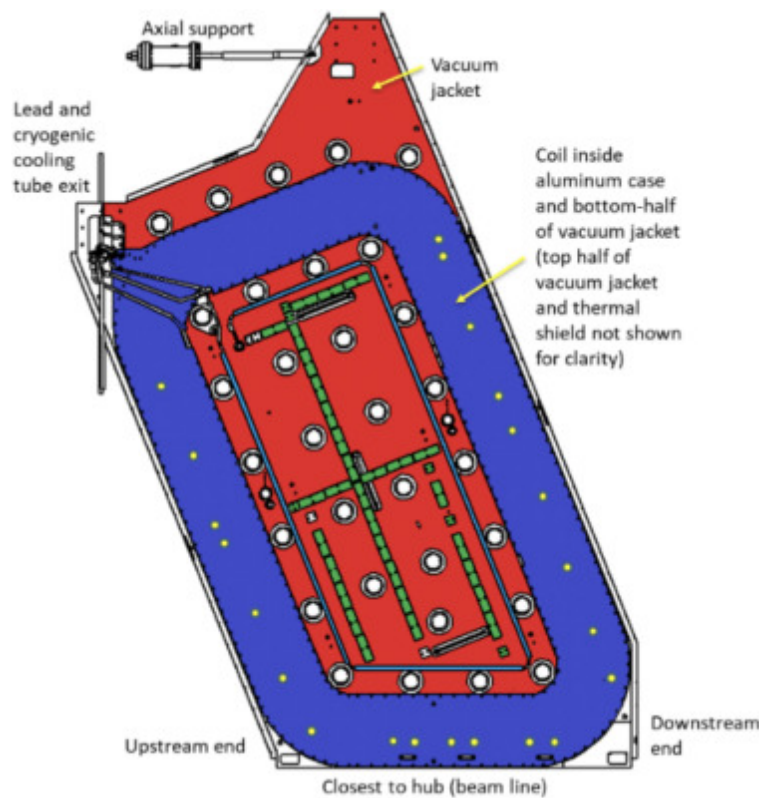


Figure 2: Magnet Line

78 This design presents several benefits as it permits a robust magnetic field within a com-
79 pact volume. The potting material offers structural support for the superconducting coils,
80 and the aluminum case aids in conducting heat away from the coils, which is crucial for

81 preserving their superconductivity.

82 1.3 Solenoid

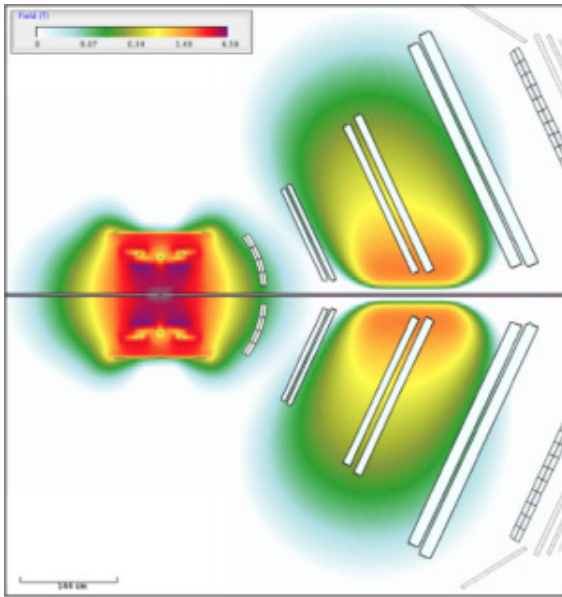


Figure 3: Magnetic fields due to the solenoid and torus



Figure 4: CLAS12 Magnets. Left: Solenoid. Right: Torus Magnet

83 The solenoid magnet is a self-shielded superconducting magnet positioned around the
84 beam trajectory, designed to produce a magnetic field along the beam direction. The design
85 is motivated by the physics requirement to safeguard the detectors in CLAS12 from scattered
86 electrons due to the Møller effect and to ensure a uniform field.

87 The magnet is composed of four cylindrical coils arranged in pairs at varying radial
88 distances along the beamline. The fifth coil is situated outside the four inner coils and
89 creates a magnetic field in the opposite direction of the inner coils' field, thereby functioning
90 as an active magnetic shield. At peak current, the solenoid produces a magnetic field of 5
91 Tesla at its core.

92 2 Beam

93 2.1 Beam Introduction

94 The CLAS12 physics program at Jefferson Lab’s Hall B relies on a versatile electron beam
95 to probe targets ranging from simple hydrogen to complex lead nuclei. To maximize the
96 scientific output, the beam needs to be precisely controlled and monitored in terms of energy,
97 current, size, and polarization. We address the challenges posed by the detector’s large
98 acceptance and proximity to the target, drawing on successful solutions implemented in
99 previous Hall B experiments. We present the key modifications made to the existing beam
100 line and the additional components introduced to achieve high-quality beams and enable
101 CLAS12 to operate at its design luminosity.[2]

102 2.2 Beam Parameters

103 The Hall B beamline is meticulously designed to meet the specific needs of CLAS12 exper-
104 iments. It offers the necessary tools to maintain precise control and ensure the safety of
105 personnel and equipment during operation.

106 The following table will detail the key parameters crucial for conducting experiments
107 with CLAS12.

Parameter	Requirement	Unit
Beam energies	≤ 11	GeV
Beam currents	< 500	nA
Current instability	≈ 10	%
Accuracy of current measurement	≈ 1	%
Beam widths (σ_x, σ_y)	< 300	μm
Position stability	< 200	μm
Divergence	< 100	μrad
Beam halo ($> 5 \sigma_x$)	$< 10^{-4}$	
Beam polarization	> 80	%
Accuracy of polarization	< 3	%

Table 1: Specifications of Beam Parameters

108 2.3 Hall B Segments

109 The Hall B beamline acts as the delivery system for CLAS12 experiments. It consists of two
110 main sections: the 2C line and the 2H line.

111 The 2C line, nicknamed the "beam switch yard delivery line," transports the beam ex-
112 tracted from CEBAF to the upstream end of the Hall B experimental hall. The 2H line then
113 takes over, guiding the beam to its destination - the beam dump located in the downstream
114 tunnel.

115 For high-energy CLAS12 operations, some upgrades were made to the beamline. The 2C
116 line now houses a new resident: the Møller polarimeter. This instrument precisely measures
117 the beam's longitudinal polarization, a key factor in certain experiments. Additionally, an
118 intermediate beam dump was installed just before the hall to provide more control over the
119 beam.[3]

120 Downstream in the 2H line, another modification was made. A cryogenic target was
121 introduced, capable of withstanding the intense beams used in CLAS12 experiments. Further
122 along the line, a tungsten shield was placed inside the CLAS12 torus magnet bore to absorb
123 any stray particles.

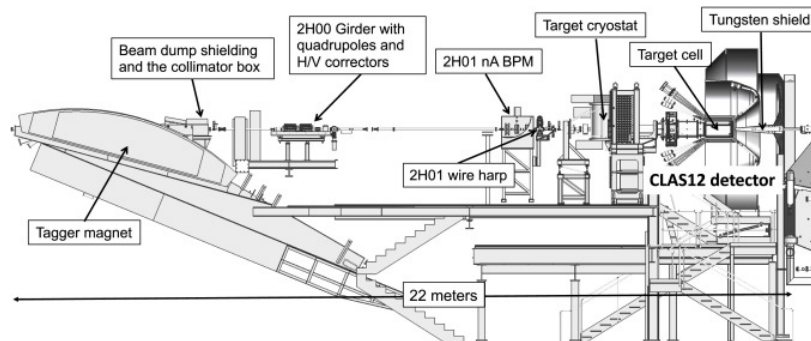


Figure 5: Beamline.

124 2.4 Instrumentation

125 The instrumentation of the beamline comprises beam current monitors, beam viewers, col-
126 limators, beam profile scanners, and beam halo monitors. Beam current monitors assist in
127 regulating the energy flow to the beam. The beam viewer enables us to see the beam, as it is

128 not visible to the naked eye. Collimators limit the beam's direction. The beam profile scan-
 129 ner manages its intensity, and halo monitors keep track of the particles that have deviated
 130 slightly.

131 On the left of Fig. 31 5, the tagger magnet remains unenergized during production data
 132 acquisition. When activated, the yoke of this magnet functions as a beam dump, utilized
 133 during beam adjustment prior to directing the beam onto the Hall B production target.
 134 Furthermore, it is employed during specific runs, such as polarization measurements in the
 135 upstream beamline, to prevent sensitive CLAS12 detectors from exposure to high background
 136 loads.

137 The status of the electron beam and every diagnostic element, including the beamline
 138 vacuum and superconducting magnets, was displayed on a monitor accessible to the shift
 139 team. More details in Fig 6.

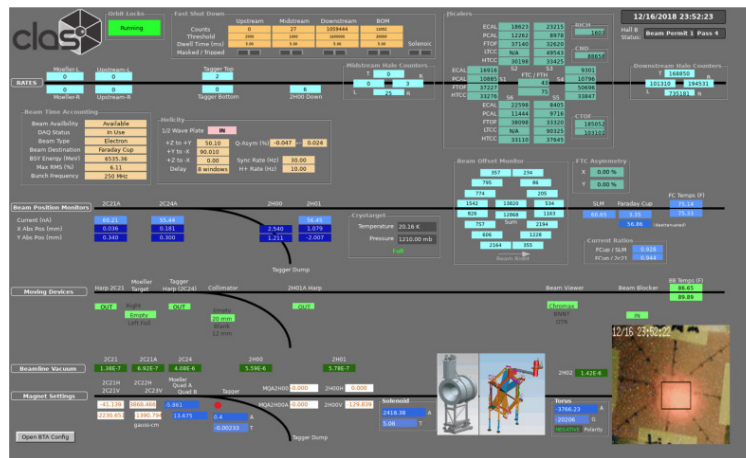


Figure 6: Beamline monitoring systems

3 Data Acquisition Simulation

Understanding the intricacies of nuclear and hadronic reactions within CLAS12 hinges on a powerful interplay between data acquisition and simulation techniques. A complex detector system, comprised of multiple subsystems like drift chambers, scintillators, and calorimeters (detailed in Section 1), captures the interactions triggered by the electron beam. This data, exceeding 100,000 channels, forms the raw material for scientific discovery.[4]

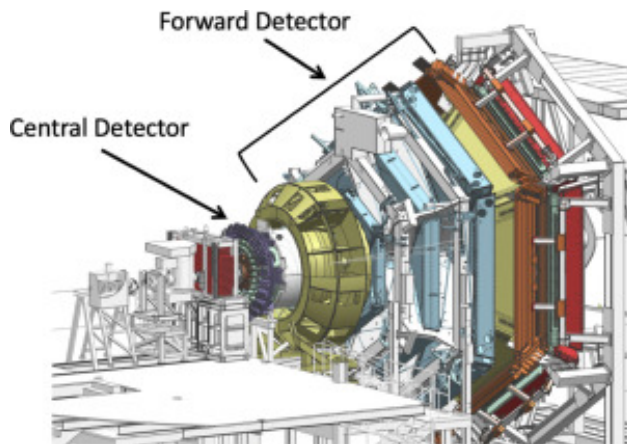


Figure 7: CLAS12.

3.1 Generated Events

The CLAS12 data-acquisition system was conceptualized as a network-based system operating on a pipeline principle. The data gathering process initiated from the front-end elements of CLAS12. These elements, despite their varying hardware and software configurations, were designed to be harmonious with the overall system. These front-end elements are known as Readout Controllers (ROCs). All ROC components were linked via TCP sockets over Ethernet to the Event Builder component. The TCP protocol guarantees that the data-sending component is linked to the data-receiving component before data transmission, unlike UCP, which transmits data before confirming the receiver's connection.

The Event Builder (EB) is a multithreaded C application that operates on a multicore Linux server. Once the events were gathered, they were forwarded to the Event Transfer (ET) system. This multithreaded C application is engineered to facilitate data-processing

158 programs to filter, monitor, and store data in shared memory. Typically, EB is executed on
159 the same server as ET, but it can also be spread across a series of ET servers. The final
160 component in this series is the Event Recorder (ER), which receives data from the ET and
161 logs them on a disk.

162 **3.2 Reconstructed Events**

163 The Event Reconstruction software, which aids in event simulation, was conceived and culti-
164 vated within the CLARA framework. CLARA assists in designing and developing scientific
165 data processing applications without the need for actual coding.

166 The data reader functionalities extract the deciphered detector data archived in the
167 database. Each deciphered detector hit entry forms a database entry, containing identifiers
168 for detector elements and data of the detector, such as current, signal and time tailored to
169 the particular system. Comparable database frameworks are produced during the decoding
170 phase for different parameters necessary for simulation. It's at this juncture that we deploy
171 reconstruction algorithms to populate these databases.

172 CLARA acts as the engine for reconstructing particle interactions within our experiment.
173 It takes data from existing databases as its raw material and generates new databases con-
174 taining reconstructed information. These outputs are then fed back into the reconstruction
175 process, leading to a more complete picture of the event.

176 The reconstruction process itself mirrors a larger framework called CLAS12. It starts
177 with reconstructing individual particles within specific detector systems, like the Central
178 and Forward Detectors. This initial step, called hit-based tracking, relies on the spatial in-
179 formation from sensor readings. Simultaneously, other detectors process data to reconstruct
180 the energy and timing of particle interactions.

181 The Event Builder service then takes center stage. It matches the reconstructed tracks
182 with energy and timing information, essentially creating a profile for each particle. Any
183 leftover unmatched data points could indicate the presence of neutral particles. At this
184 point, the Event Builder can pinpoint the exact moment of interaction between the beam

185 and the target, effectively setting the event’s ”start time.” With this crucial information, a
186 more refined tracking process, called time-based tracking, can be initiated.[5]

187 The improved tracks are then fed back into the Event Builder, leading to a final, com-
188 prehensive reconstruction of the entire event. Importantly, some reconstruction steps can
189 happen concurrently, while others require the completion of previous steps before they can
190 begin.

191 To streamline this complex process, we’ve developed a service-oriented architecture specif-
192 ically designed for CLAS12 data reconstruction and analysis. This framework utilizes vari-
193 ous libraries to manage data input/output, detector geometry, databases, and magnetic field
194 configurations, ensuring the smooth operation of the entire reconstruction pipeline.

195 **4 CLAS12 Experiments**

196 First, we embark on a journey with the Heavy Photon Search (HPS) experiment [6]. Here,
197 CLAS12 serves as a vital component in the hunt for a new particle, the elusive ”heavy
198 photon.”

199 Secondly, we shift focus to the realm of hadron interactions, where CLAS12 plays a key
200 role in unraveling the mysteries of beam spin asymmetries in pion production (See Figure
201 4). This experiment sheds light on the intricate dance between gluons and quarks within
202 protons [7].

203 **4.1 Heavy Photon Search**

204 **4.1.1 Objective**

205 In order for scientists to search for the dark photon in fixed target electro-production, they
206 design the Heavy Photon Search at Thomas Jefferson National Accelerator Facility. Their
207 purpose is to detect e^-e^+ decay of the dark photon.

208 For them to sense the decay a silicon vertex tracker and a fragmented electromagnetic
209 calorimeter is used to measure the bump on the data.

210 **4.1.2 Beamline Characteristics**

211 The layout consists of our target inside a vacuum chamber with our silicon vertex trackers.
 212 This is inside the center of our dipole magnets core, behind the magnets is our calorimeter.
 213 The entirety of this setup is located in a compartment of the CLAS12.

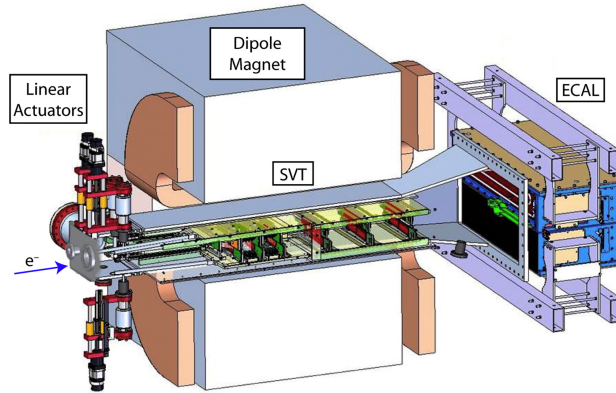


Figure 8: Cross-section of HPS setup

214 The experiment runs using beam energies that range from 1 to 6.6 GeV and beam currents
 215 fluctuating from 1 to 500nA. A table with the beam parameters is below.

Parameter	Requirement	Unit
Beam energies	$\leq 1 - 6.6$	GeV
Beam currents	< 500	nA
Current instability	≈ 5	%
Beam widths (σ_x, σ_y)	$< (300, 50)$	μm
Position stability	< 30	μm
Divergence	< 100	μrad
Beam halo ($> 5 \sigma_x$)	$< 10^{-5}$	

Table 2: Beam Parameters HPS

216 This parameters for the beam have been established by running a variety of simulations.
 217 The two segments of the Hall B are use for this experiment. The 2C segment redirects
 218 and elevates our beam to the Hall B elevation after its extracted from the Beam Switch Yard
 219 (BSY). In the 2H segment, specifically near the end of the beamline, the target is located.
 220 The target being tungsten foil. The devices located in the 2H segment are enough to monitor
 221 the beam properties.

222 **4.1.3 Performance**

223 Before starting production is necessary to evaluate the quality of the beam. So a first tuning
224 is done to correct the beam profile in the 2C segment. After, the beam is sent to the beam
225 dump on Hall B to fix the profile of the downstream beam.

226 After the beam profile is perfected the orbit lock system is displayed. This system helps
227 control and monitored the beam width to regulate the motion with respect to the target.
228 Succeeding the verification of the beam profile the HPS target is inserted. Following the
229 target, the halo counters of the FSD are inserted. This helps control the profile and width
230 because if the beam halo moves closer to the collimator walls or the detectors in the edges
231 the halo rates would increase.

232 When the beam evaluation is being done in the 2H segment for the first time a silicon
233 vertex tracker scan is executed in order to control the profile and to check the beam alignment
234 with respect to the silicon vertex tracker coordinate system. In Fig. 9 the first peak shows the
235 halo rate with respect to the SVTs coordinates. If it was not center it would be adjusted.[6]

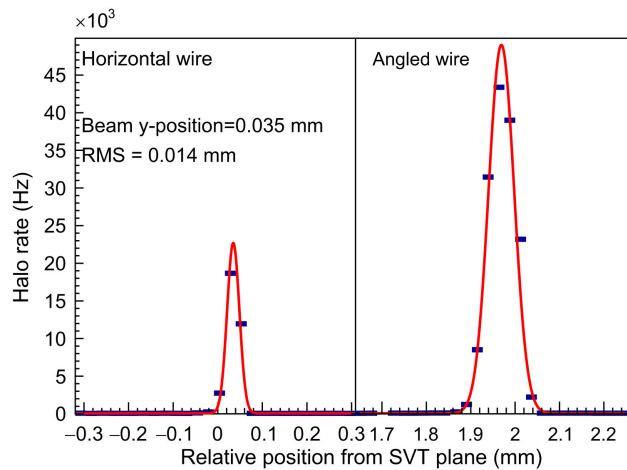


Figure 9: SVT Scan

236 **4.2 Beam Spin Asymmetries**

237 **4.2.1 Objective**

238 Parton distribution functions "PDF" and Fragmentation Functions "FF" are of high impor-
239 tance to particle physicists because they contain information of the momentum distribution

240 of the quarks and of the formation of hadrons. For the purpose of accessing them, semi-
 241 inclusive deep inelastic scattering "SIDIS" is considered. Specifically the high energy of the
 242 electron collision with the proton allows us to analyze the interaction as if it were a one-on-
 243 one encounter with a single fundamental particle within the proton. This is why the single
 244 spin asymmetries for the production of two pions was studied at JLab. [7]

245
$$e(l) + p(P) \rightarrow e'(l') + \pi^+(P_1) + \pi^-(P_2) + X$$

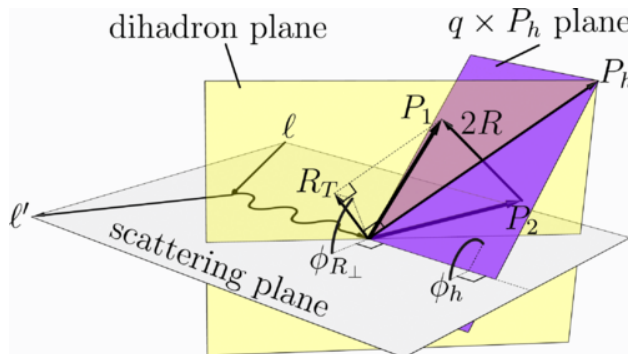


Figure 10: Reaction Plane

246 In the figure above we have the plane of our reaction. The scattering plane is mainly
 247 defined by l , which can be visualized as the incoming electron beam and the dihadron plane
 248 is created by P_1 and P_2 , the initial being the hadron momentum.

249 **4.2.2 Detectors and Data Extraction**

250 Experimental data was acquire using the spectrometer located in the CLAS12. The beam
 251 was longitudinally polarized with an energy of 10.6 GeV and with a frequency of 30 Hz it was
 252 inverted in order to diminish systematic errors. The beam was acquired through CEBAF
 253 (Continuous Electron Beam Accelerator Facility), subsequently it enter the 2C segment of
 254 Hall B where the beam profile was tuned before entering the 2H segment.

255 After the tuning was performed on the 2H segment the target, which consisted of liquid
 256 hydrogen, was placed before the CLAS12 detector.

257 For the purpose of identifying the scattered electrons and pions the Forward Detector
 258 is used. The Forward Detector is composed of a tracking subsystem which identifies the
 259 electrons and pions through a set of Cherenkov counters and drift chambers inside a toroid.

260 This method coupled with an supplementary identification by an electromagnetic calorimeter
 261 for electrons and plastic scintillators for pions ensures identification of them.

262 4.2.3 Findings

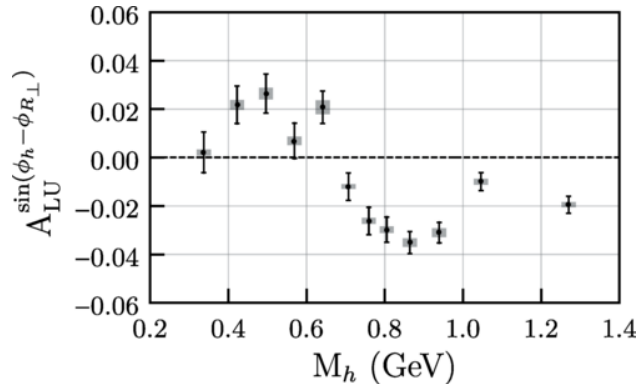


Figure 11: A_{LU} vs M_h (GeV)

263 The key finding of the experiment was the observation of a clear change in the sign of
 264 the beam spin asymmetry around the mass of the rho meson (ρ). The change in sign in
 265 Fig.11 indicates a dependence between the helicity (spin) of the fragmenting quark and the
 266 produced pions. In simpler terms, the spin orientation of the quark leaving the collision
 267 seems to influence the type of pions created. This finding provides the first experimental
 268 signal for a specific quantity called the helicity-dependent two-pion fragmentation function.
 269 This function offers valuable insights into the interactions between gluons and quarks within
 270 protons.

271 5 Conclusions

272 The CLAS12 detector stands as a remarkable feat of human ingenuity, a monument to our
 273 relentless quest to crack the code of the subatomic world. Its innovative design, a marriage
 274 of powerful toroidal and solenoid magnets, shatters the boundaries of particle detection.
 275 CLAS12 allows us to capture a stunning array of charged and neutral particles across a
 276 significant chunk of this hidden universe. This comprehensive data collection forms the
 277 foundation for groundbreaking discoveries in nuclear and hadronic physics, each finding a

278 stepping stone towards a deeper understanding of the cosmos.

279 The success of CLAS12 experiments is a delicate dance between cutting-edge technology
280 and meticulous control. The precisely tuned beamline components ensure the delivery of
281 electron beams with the desired energy, current, size, and polarization. Sophisticated instru-
282 ments within the beamline act as guardians, continuously monitoring and maintaining these
283 parameters. Any deviation could introduce noise into the delicate symphony of subatomic
284 interactions, potentially masking the data.

285 The data acquisition and simulation systems of CLAS12 act as the bridge between the
286 captured data and the secrets it holds. The complex detector system meticulously captures
287 the fleeting interactions, transforming them into raw data. This data then undergoes a
288 sophisticated transformation through advanced software tools. Reconstruction algorithms
289 breathe life into the data, piecing together the fragments of the collision events. Through
290 this process, scientists gain unparalleled access to the fundamental forces and building blocks
291 of matter, the quarks and gluons that reside within protons and neutrons.

292 The examples presented here; the Heavy Photon Search (HPS) and the study of Beam
293 Spin Asymmetries in pion production showcase just a glimpse of CLAS12's immense poten-
294 tial. The HPS experiment allows scientists to peer beyond the Standard Model of particle
295 physics, searching for entirely new particles that could hold the key to a deeper understand-
296 ing of the universe. Meanwhile, the studies of beam spin asymmetries shed light on the
297 complex internal dynamics of protons, revealing the intricate interplay between gluons and
298 quarks.

299 Looking ahead, CLAS12 remains poised to play a pivotal role in the ongoing exploration
300 of the subatomic world. As scientists continue to push the boundaries of this realm, CLAS12
301 stands ready to capture the fleeting glimpses of these interactions, offering invaluable data
302 for groundbreaking discoveries in the years to come.

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